

# Section 16.7 Stokes' Theorem

Recall that for a v.f.  $F = \langle M, N \rangle$  on the plane we have

$$\text{div } F = \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} = \begin{pmatrix} \text{measures compression} \\ \text{or expansion at } (x, y) \end{pmatrix}$$

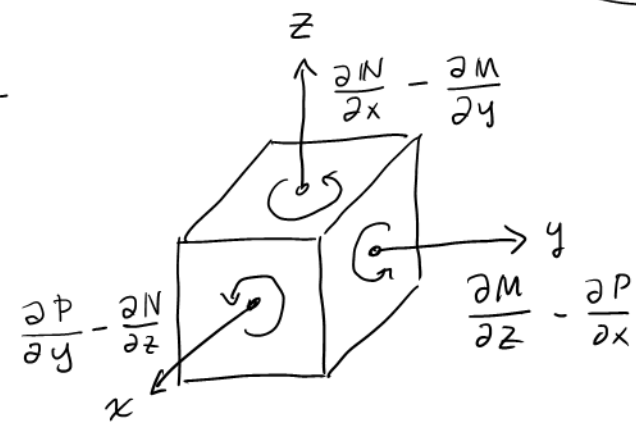
$$\text{curl } F = \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} = \begin{pmatrix} \text{measures counterclockwise} \\ \text{circulation at } (x, y) \end{pmatrix}$$

Now lets adapt all this to 3-D. Let  $F = \langle M, N, P \rangle$  and think of it as representing the velocity of a fluid or gas flowing in space.

## Divergence

$$\text{div } F = \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} + \frac{\partial P}{\partial z} = \underbrace{\left\langle \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right\rangle}_{\Delta} \cdot \underbrace{\langle M, N, P \rangle}_F = \nabla \cdot F$$

## Curl

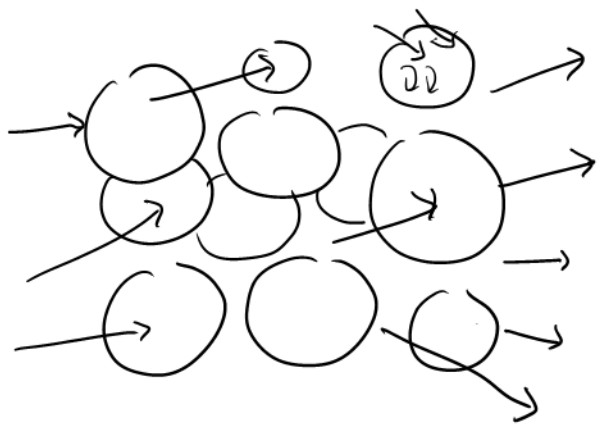


special notation

$$\text{curl } F = \left\langle \frac{\partial P}{\partial y} - \frac{\partial N}{\partial z}, \frac{\partial M}{\partial z} - \frac{\partial P}{\partial x}, \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right\rangle = \begin{pmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ M & N & P \end{pmatrix} = \nabla \times F$$

$= \left( \begin{array}{l} \text{Axis of spin} \\ \text{in v.f. } F \text{ at} \\ \text{point } (x, y, z) \end{array} \right)$ 
 $\left. \begin{array}{l} |\text{curl } F| = \\ \text{speed of} \\ \text{spin} \end{array} \right)$

Imagine v.f.  $F$  moving little balls through space. As they flow through space, they also may spin.

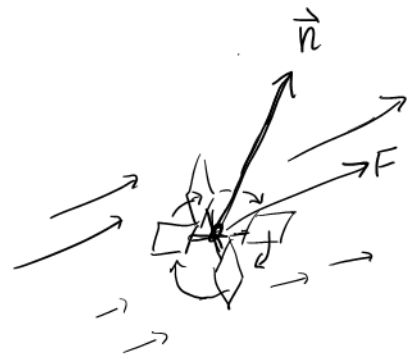


$$F(x, y, z) = \text{velocity of ball at } (x, y, z)$$

$$\text{div } F = \begin{pmatrix} \text{Compression of balls} \\ \text{at } (x, y, z) \end{pmatrix}$$

$$\text{curl } F = \begin{pmatrix} \text{spin of ball} \\ \text{at } (x, y, z) \end{pmatrix}$$

Given a unit vector  $\vec{n}$ , imagine that it has a little paddle wheel at its end. The v.f.  $F$  would cause it to spin.



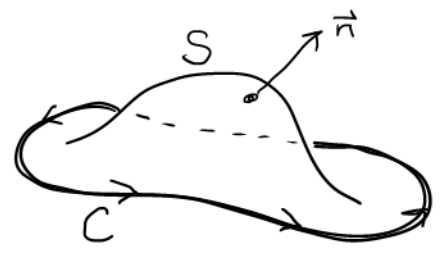
$$(\text{curl } F) \cdot \vec{n} = \nabla \times F \cdot \vec{n} = \begin{cases} \text{a measure of the spin of } \vec{n}. \\ \text{positive: } \begin{array}{c} \text{---} \rightarrow \vec{n} \\ \curvearrowright \end{array} \\ \text{negative: } \begin{array}{c} \text{---} \rightarrow \vec{n} \\ \curvearrowleft \end{array} \\ \text{Zero - no spin} \\ \text{bigger = faster spin} \end{cases}$$

With all this in mind, we can state and understand the main result of this section - Stokes' Theorem.

### Stokes' Theorem

Suppose  $F = \langle M, N, P \rangle$  is a vector field in space and  $S$  is a surface with normal  $\vec{n}$  and boundary  $C$ , a curve  $\vec{r}(t)$  traversed counterclockwise (looking down  $\vec{n}$ ). Then:

$$\oint_C F \cdot d\vec{r} = \iint_S \nabla \times F \cdot \vec{n} \, d\sigma$$

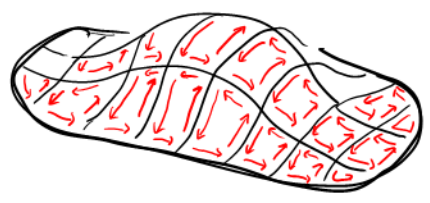


i.e. (circulation around  $C$ ) =  $\iint_S$  (circulation at point  $(x, y, z)$  on surface)  $d\sigma$

$\nabla \times F \cdot \vec{n} = (\text{curl } F) \cdot \vec{n} =$  "flux of curl"

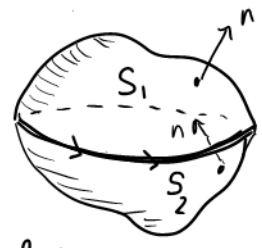
You can understand this intuitively by dividing  $S$  up into small rectangles

$$\iint_S \nabla F \cdot n \, d\sigma = \lim_{|P| \rightarrow 0} \sum \nabla f \cdot n \, \Delta \sigma_k$$



the contributions to circulation along adjacent rectangles cancel. The only contribution is along the boundary  $C$ ; it is  $\oint_C F \cdot d\vec{r}$ .

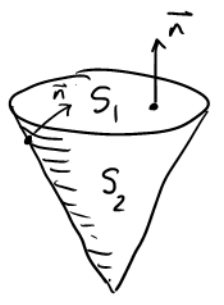
Notice that the theorem implies that if surfaces  $S_1$  and  $S_2$  share a common boundary, then



$$\iint_{S_1} \nabla \times F \cdot \vec{n} \, d\sigma = \iint_{S_2} \nabla \times F \cdot \vec{n} \, d\sigma$$

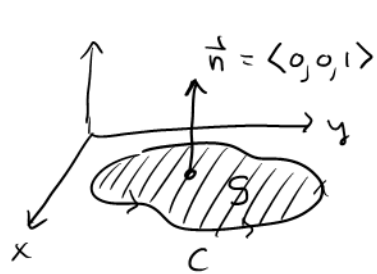
because both sides equal  $\oint_C F \cdot d\vec{r}$

For instance suppose  $S_1$  is top of cone and  $S_2$  is side. Then



$$\iint_{S_1} \nabla \times F \cdot \vec{n} \, d\sigma = \iint_{S_2} \nabla \times F \cdot \vec{n} \, d\sigma$$

Notice also what happens if  $S$  is a flat surface on the  $xy$ -plane



$$\oint_C \mathbf{F} \cdot d\mathbf{r} = \iint_S (\nabla \times \mathbf{F}) \cdot \langle 0, 0, 1 \rangle d\sigma$$

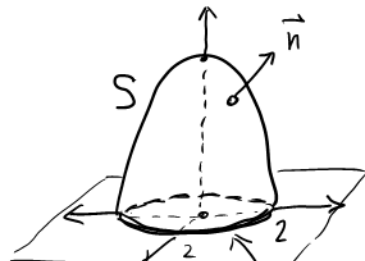
$$\Rightarrow \oint_C \mathbf{F} \cdot T ds = \iint_S \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$$

In this case, Stokes' theorem reduces to Green's Theorem. In other words, Stokes theorem is a generalization to Green's Theorem.

Example  $\mathbf{F}(x, y, z) = \langle y, -x^2, 2z^2 \rangle$

$S$  is the part of  $4 - x^2 - y^2$  above  $xy$ -plane

Verify that Stokes' theorem holds.



Left hand side:

$$\oint_C \mathbf{F} \cdot d\mathbf{r} = \oint_C \mathbf{F} \cdot \frac{d\mathbf{r}}{dt} dt$$

$$= \int_0^{2\pi} \langle 2\sin t, -(2\cos t)^2, 2 \cdot 0^2 \rangle \cdot \langle -2\sin t, 2\cos t, 0 \rangle dt$$

$$= \int_0^{2\pi} (-4\sin^2 t - 8\cos^3 t) dt = \int_0^{2\pi} \left( -4 \frac{1 - \cos(2t)}{2} - 8\cos^2 t \cos t \right) dt$$

$$= \int_0^{2\pi} -2(1 - \cos 2t) - 8(1 - \sin^2 t) \cos t dt$$

$$= \int_0^{2\pi} \left( -2 + \cos 2t - 8\cos t + 8(\sin t)^2 \cos t \right) dt$$

$$= \left[ -2t + \frac{1}{2} \sin 2t - 8\sin t + 8 \frac{\sin^3 t}{3} \right]_0^{2\pi}$$

$$= \left( -4\pi + \frac{1}{2} \sin 4\pi - 8\sin 2\pi + 8 \frac{\sin^3 2\pi}{3} \right) - \left( -2 \cdot 0 + \frac{1}{2} \sin 0 - 8\sin 0 + 8 \frac{\sin^3 0}{3} \right)$$

$$= \boxed{-4\pi}$$

Base circle  $C$   
of radius 2  
is boundary of  $S$   
 $\vec{r}(t) = \langle 2\cos t, 2\sin t, 0 \rangle$   
 $0 \leq t \leq 2\pi$

$u = \sin t$   
 $du = \cos t dt$

Now let's compute the right-hand side.

$$\iint_S \nabla \times F \cdot n \, d\sigma$$

First,  $\nabla \times F = \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y & -x^2 & 2z^2 \end{vmatrix} = \langle 0, 0, -2x-1 \rangle$  hold that thought

Next, we parameterize the surface S as:

$$\vec{r}(u,v) = \langle u \cos v, u \sin v, 4 - (u \cos v)^2 - (u \sin v)^2 \rangle \left. \vphantom{\vec{r}(u,v)} \right\} 0 \leq u \leq 2, 0 \leq v \leq 2\pi$$

$$= \langle u \cos v, u \sin v, 4 - u^2 \rangle$$

$$\vec{r}_u = \langle \cos v, \sin v, -2u \rangle$$

$$\vec{r}_v = \langle -u \sin v, u \cos v, 0 \rangle$$

$$\vec{r}_u \times \vec{r}_v = \begin{vmatrix} i & j & k \\ \cos v & \sin v & -2u \\ -u \sin v & u \cos v & 0 \end{vmatrix} = \langle 2u^2 \cos v, -2u^2 \sin v, u \rangle$$

not needed

$$|\vec{r}_u \times \vec{r}_v| = \sqrt{(2u^2 \cos v)^2 + (-2u^2 \sin v)^2 + u^2} = \sqrt{4u^4 \cos^2 v + 4u^4 \sin^2 v + u^2}$$

$$= \sqrt{4u^4 (\cos^2 v + \sin^2 v) + u^2} = \sqrt{4u^4 + u^2} = u \sqrt{4u^2 + 1}$$

Now:  $\iint_S \nabla \times F \cdot n \, d\sigma = \int_0^{2\pi} \int_0^2 \underbrace{\langle 0, 0, -2x-1 \rangle}_{\nabla \times F} \cdot \underbrace{\frac{\vec{r}_u \times \vec{r}_v}{|\vec{r}_u \times \vec{r}_v|}}_{\vec{n}} \underbrace{|\vec{r}_u \times \vec{r}_v|}_{d\sigma} \, du \, dv$

$$= \int_0^{2\pi} \int_0^2 \langle 0, 0, -2x-1 \rangle \cdot (\vec{r}_u \times \vec{r}_v) \, du \, dv$$

$$= \int_0^{2\pi} \int_0^2 \langle 0, 0, -2u \cos v - 1 \rangle \cdot \langle 2u^2 \cos v, -2u^2 \sin v, u \rangle \, du \, dv$$

$$= \int_0^{2\pi} \int_0^2 (-2u^2 \cos v - u) \, du \, dv = \int_0^{2\pi} \left[ -\frac{2}{3} u^3 \cos v - \frac{u^2}{2} \right]_0^2 \, dv = \int_0^{2\pi} \left( -\frac{16}{3} \cos v - 2 \right) \, dv$$

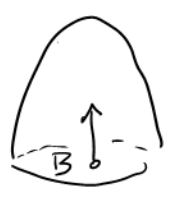
$$= \left[ -\frac{16}{3} \sin v - 2v \right]_0^{2\pi} = \left( -\frac{16}{3} \sin 2\pi - 2(2\pi) \right) - \left( -\frac{16}{3} \sin 0 - 2 \cdot 0 \right) = \boxed{-4\pi}$$

Notice we could also integrate over the circular base & get the same answer:

$$\iint_B \nabla \times F \cdot n \, d\sigma = \iint_B \langle 0, 0, -2x-1 \rangle \cdot \langle 0, 0, 1 \rangle \, dA$$

switch to polar

$$= \iint_B (-2x-1) \, dA = \int_0^{2\pi} \int_0^2 (-2r \cos \theta - 1) r \, dr \, d\theta = \boxed{-4\pi}$$



Besides having many physical applications, Stokes theorem can also help evaluate a line- or surface-integral, if the other side is easier to do.